

# Lectures on Lévy Processes and Stochastic Calculus, Braunschweig, Lecture 2: Lévy Processes

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We say that  $X$  is a *Lévy process* if

(L1) Each  $X(0) = 0$  (a.s),

(L2)  $X$  has independent and stationary increments,

(L3)  $X$  is *stochastically continuous* i.e. for all  $a > 0$  and for all  $s \geq 0$ ,

$$\lim_{t \rightarrow s} P(|X(t) - X(s)| > a) = 0.$$

Note that in the presence of (L1) and (L2), (L3) is equivalent to the condition

$$\lim_{t \downarrow 0} P(|X(t)| > a) = 0.$$

## Definition: Lévy Process

Let  $X = (X(t), t \geq 0)$  be a stochastic process defined on a probability space  $(\Omega, \mathcal{F}, P)$ .

We say that it has *independent increments* if for each  $n \in \mathbb{N}$  and each

$0 \leq t_1 < t_2 < \dots < t_{n+1} < \infty$ , the random variables  $(X(t_{j+1}) - X(t_j), 1 \leq j \leq n)$  are independent

and it has *stationary increments* if each

$$X(t_{j+1}) - X(t_j) \stackrel{d}{=} X(t_{j+1} - t_j) - X(0).$$

The *sample paths* of a process are the maps  $t \rightarrow X(t)(\omega)$  from  $\mathbb{R}^+$  to  $\mathbb{R}^d$ , for each  $\omega \in \Omega$ .

We are now going to explore the relationship between Lévy processes and infinite divisibility.

### Theorem

If  $X$  is a Lévy process, then  $X(t)$  is infinitely divisible for each  $t \geq 0$ .

*Proof.* For each  $n \in \mathbb{N}$ , we can write

$$X(t) = Y_1^{(n)}(t) + \dots + Y_n^{(n)}(t)$$

where each  $Y_k^{(n)}(t) = X(\frac{kt}{n}) - X(\frac{(k-1)t}{n})$ .

The  $Y_k^{(n)}(t)$ 's are i.i.d. by (L2). □

From Lecture 1 we can write  $\phi_{X(t)}(u) = e^{\eta(t,u)}$  for each  $t \geq 0, u \in \mathbb{R}^d$ , where each  $\eta(t, \cdot)$  is a Lévy symbol.

### Theorem

If  $X$  is a Lévy process, then

$$\phi_{X(t)}(u) = e^{t\eta(u)},$$

for each  $u \in \mathbb{R}^d, t \geq 0$ , where  $\eta$  is the Lévy symbol of  $X(1)$ .

*Proof.* Suppose that  $X$  is a Lévy process and for each  $u \in \mathbb{R}^d, t \geq 0$ , define  $\phi_u(t) = \phi_{X(t)}(u)$  then by (L2) we have for all  $s \geq 0$ ,

$$\begin{aligned} \phi_u(t+s) &= \mathbb{E}(e^{i(u, X(t+s))}) \\ &= \mathbb{E}(e^{i(u, X(t+s)-X(s))} e^{i(u, X(s))}) \\ &= \mathbb{E}(e^{i(u, X(t+s)-X(s))}) \mathbb{E}(e^{i(u, X(s))}) \\ &= \phi_u(t) \phi_u(s) \dots (i) \end{aligned}$$

X

Now  $\phi_u(0) = 1 \dots (ii)$  by (L1), and the map  $t \rightarrow \phi_u(t)$  is continuous. However the unique continuous solution to (i) and (ii) is given by  $\phi_u(t) = e^{t\alpha(u)}$ , where  $\alpha : \mathbb{R}^d \rightarrow \mathbb{C}$ . Now by Theorem 1,  $X(1)$  is infinitely divisible, hence  $\alpha$  is a Lévy symbol and the result follows. □

We now have the Lévy-Khinchine formula for a Lévy process  $X = (X(t), t \geq 0)$ :-

$$\begin{aligned} \mathbb{E}(e^{i(u, X(t))}) &= \exp\left\{ t \left[ i(b, u) - \frac{1}{2}(u, Au) \right. \right. \\ &\quad \left. \left. + \int_{\mathbb{R}^d - \{0\}} (e^{i(u, y)} - 1 - i(u, y) \mathbf{1}_{\tilde{B}}(y)) \nu(dy) \right] \right\}, (2.1) \end{aligned}$$

for each  $t \geq 0, u \in \mathbb{R}^d$ , where  $(b, A, \nu)$  are the characteristics of  $X(1)$ . We will define the Lévy symbol and the characteristics of a Lévy process  $X$  to be those of the random variable  $X(1)$ . We will sometimes write the former as  $\eta_X$  when we want to emphasise that it belongs to the process  $X$ .

Let  $p_t$  be the law of  $X(t)$ , for each  $t \geq 0$ . By (L2), we have for all  $s, t \geq 0$  that:

$$p_{t+s} = p_t * p_s.$$

By (L3), we have  $p_t \xrightarrow{w} \delta_0$  as  $t \rightarrow 0$ , i.e.  $\lim_{t \rightarrow 0} \int f(x) p_t(dx) = f(0)$ .

$(p_t, t \geq 0)$  is a *weakly continuous convolution semigroup of probability measures* on  $\mathbb{R}^d$ .

Conversely, given any such semigroup, we can always construct a Lévy process on path space via Kolmogorov's construction.

Informally, we have the following asymptotic relationship between the law of a Lévy process and its Lévy measure:

$$\nu = \lim_{t \downarrow 0} \frac{p_t}{t}.$$

More precisely

$$\lim_{t \downarrow 0} \frac{1}{t} \int_{\mathbb{R}^d} f(x) p_t(dx) = \int_{\mathbb{R}^d} f(x) \nu(dx), \quad (2.2)$$

for bounded, continuous functions  $f$  which vanish in some neighborhood of the origin.

## Examples of Lévy Processes

### Example 1, Brownian Motion and Gaussian Processes

A (standard) Brownian motion in  $\mathbb{R}^d$  is a Lévy process

$B = (B(t), t \geq 0)$  for which

(B1)  $B(t) \sim N(0, tI)$  for each  $t \geq 0$ ,

(B2)  $B$  has continuous sample paths.

It follows immediately from (B1) that if  $B$  is a standard Brownian motion, then its characteristic function is given by

$$\phi_{B(t)}(u) = \exp\left\{-\frac{1}{2}t|u|^2\right\},$$

for each  $u \in \mathbb{R}^d, t \geq 0$ .

We introduce the marginal processes  $B_i = (B_i(t), t \geq 0)$  where each  $B_i(t)$  is the  $i$ th component of  $B(t)$ , then it is not difficult to verify that the  $B_i$ 's are mutually independent Brownian motions in  $\mathbb{R}$ . We will call these *one-dimensional Brownian motions* in the sequel.

Brownian motion has been the most intensively studied Lévy process. In the early years of the twentieth century, it was introduced as a model for the physical phenomenon of Brownian motion by Einstein and Smoluchowski and as a description of the dynamical evolution of stock prices by Bachelier.

The theory was placed on a rigorous mathematical basis by Norbert Wiener in the 1920's.

We could try to use the Kolmogorov existence theorem to construct one-dimensional Brownian motion from the following prescription on cylinder sets of the form

$I_{t_1, \dots, t_n}^H = \{\omega \in \Omega; \omega(t_1) \in [a_1, b_1], \dots, \omega(t_n) \in [a_n, b_n]\}$  where  $H = [a_1, b_1] \times \dots \times [a_n, b_n]$  and we have taken  $\Omega$  to be the set of all mappings from  $\mathbb{R}^+$  to  $\mathbb{R}$ :

$$P(I_{t_1, \dots, t_n}^H) = \int_H \frac{1}{(2\pi)^{\frac{n}{2}} \sqrt{t_1(t_2 - t_1) \dots (t_n - t_{n-1})}} \exp\left(-\frac{1}{2} \left( \frac{x_1^2}{t_1} + \frac{(x_2 - x_1)^2}{t_2 - t_1} + \dots + \frac{(x_n - x_{n-1})^2}{t_n - t_{n-1}} \right)\right) dx_1 \dots dx_n.$$

However there there is then no guarantee that the paths are continuous.

The literature contains a number of ingenious methods for constructing Brownian motion. One of the most delightful of these (originally due to Paley and Wiener) obtains this, in the case  $d = 1$ , as a random Fourier series for  $0 \leq t \leq 1$ :

$$B(t) = \frac{\sqrt{2}}{\pi} \sum_{n=0}^{\infty} \frac{\sin(\pi t(n + \frac{1}{2}))}{n + \frac{1}{2}} \xi(n),$$

for each  $t \geq 0$ , where  $(\xi(n), n \in \mathbb{N})$  is a sequence of i.i.d.  $N(0, 1)$  random variables.

We list a number of useful properties of Brownian motion in the case  $d = 1$ .

- Brownian motion is locally Hölder continuous with exponent  $\alpha$  for every  $0 < \alpha < \frac{1}{2}$  i.e. for every  $T > 0, \omega \in \Omega$  there exists  $K = K(T, \omega)$  such that

$$|B(t)(\omega) - B(s)(\omega)| \leq K|t - s|^\alpha,$$

for all  $0 \leq s < t \leq T$ .

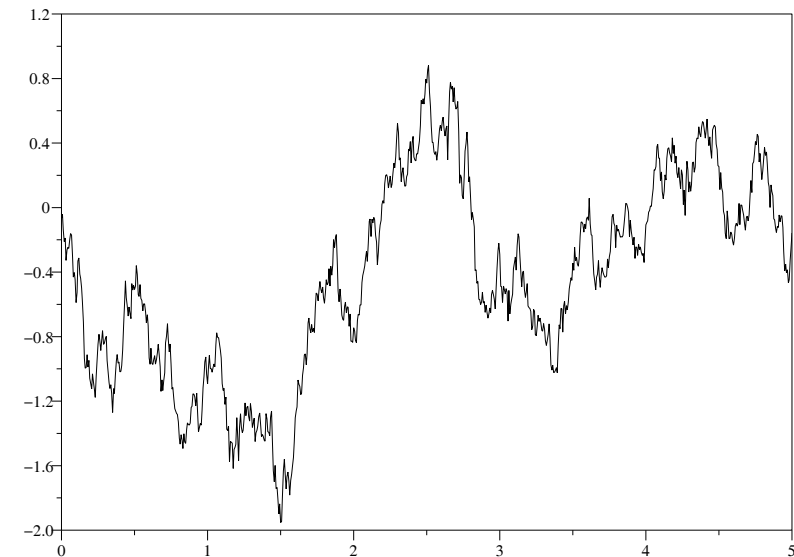
- The sample paths  $t \rightarrow B(t)(\omega)$  are almost surely nowhere differentiable.

- For any sequence,  $(t_n, n \in \mathbb{N})$  in  $\mathbb{R}^+$  with  $t_n \uparrow \infty$ ,

$$\liminf_{n \rightarrow \infty} B(t_n) = -\infty \text{ a.s. } \quad \limsup_{n \rightarrow \infty} B(t_n) = \infty \text{ a.s.}$$

- The law of the iterated logarithm:-

$$P\left(\limsup_{t \downarrow 0} \frac{B(t)}{(2t \log(\log(\frac{1}{t})))^{\frac{1}{2}}} = 1\right) = 1.$$



Simulation of standard Brownian motion

Let  $A$  be a non-negative symmetric  $d \times d$  matrix and let  $\sigma$  be a square root of  $A$  so that  $\sigma$  is a  $d \times m$  matrix for which  $\sigma\sigma^T = A$ . Now let  $b \in \mathbb{R}^d$  and let  $B$  be a Brownian motion in  $\mathbb{R}^m$ . We construct a process  $C = (C(t), t \geq 0)$  in  $\mathbb{R}^d$  by

$$C(t) = bt + \sigma B(t), \quad (2.3)$$

then  $C$  is a Lévy process with each  $C(t) \sim N(bt, tA)$ . It is not difficult to verify that  $C$  is also a Gaussian process, i.e. all its finite dimensional distributions are Gaussian. It is sometimes called *Brownian motion with drift*. The Lévy symbol of  $C$  is

$$\eta_C(u) = i(b, u) - \frac{1}{2}(u, Au).$$

In fact a Lévy process has continuous sample paths if and only if it is of the form (2.3).

## Example 2 - The Poisson Process

The Poisson process of intensity  $\lambda > 0$  is a Lévy process  $N$  taking values in  $\mathbb{N} \cup \{0\}$  wherein each  $N(t) \sim \pi(\lambda t)$  so we have

$$P(N(t) = n) = \frac{(\lambda t)^n}{n!} e^{-\lambda t},$$

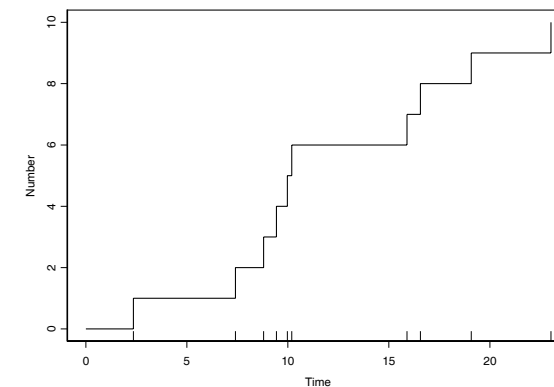
for each  $n = 0, 1, 2, \dots$

The Poisson process is widely used in applications and there is a wealth of literature concerning it and its generalisations.

We define non-negative random variables  $(T_n, \mathbb{N} \cup \{0\})$  (usually called waiting times) by  $T_0 = 0$  and for  $n \in \mathbb{N}$ ,

$$T_n = \inf\{t \geq 0; N(t) = n\},$$

then it is well known that the  $T_n$ 's are gamma distributed. Moreover, the inter-arrival times  $T_n - T_{n-1}$  for  $n \in \mathbb{N}$  are i.i.d. and each has exponential distribution with mean  $\frac{1}{\lambda}$ . The sample paths of  $N$  are clearly piecewise constant with "jump" discontinuities of size 1 at each of the random times  $(T_n, n \in \mathbb{N})$ .



Simulation of a Poisson process ( $\lambda = 0.5$ )

For later work it is useful to introduce the *compensated Poisson process*  $\tilde{N} = (\tilde{N}(t), t \geq 0)$  where each  $\tilde{N}(t) = N(t) - \lambda t$ . Note that  $\mathbb{E}(\tilde{N}(t)) = 0$  and  $\mathbb{E}(\tilde{N}(t)^2) = \lambda t$  for each  $t \geq 0$ .

### Example 3 - The Compound Poisson Process

Let  $(Z(n), n \in \mathbb{N})$  be a sequence of i.i.d. random variables taking values in  $\mathbb{R}^d$  with common law  $\mu_Z$  and let  $N$  be a Poisson process of intensity  $\lambda$  which is independent of all the  $Z(n)$ 's. The *compound Poisson process*  $Y$  is defined as follows:-

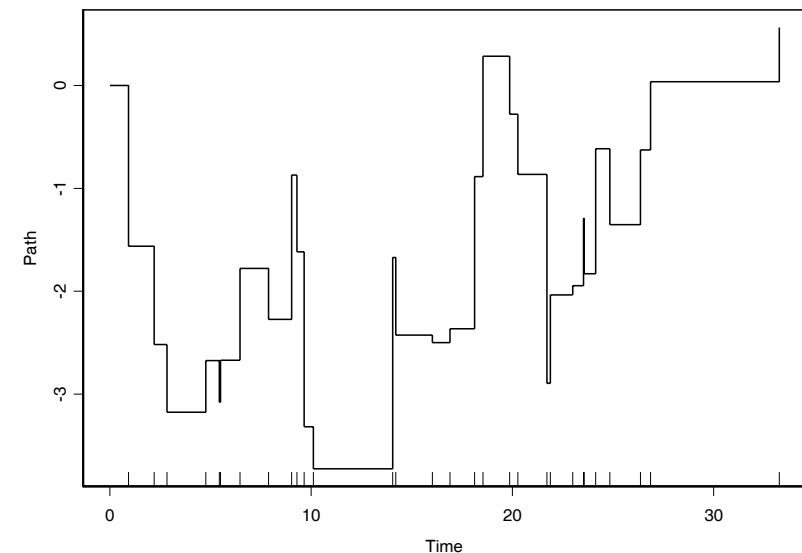
$$Y(t) := \begin{cases} 0 & \text{if } N(t) = 0 \\ Z(1) + \dots + Z(N(t)) & \text{if } N(t) > 0, \end{cases}$$

for each  $t \geq 0$ ,  
so each  $Y(t) \sim \pi(\lambda t, \mu_Z)$ .

From the work of Lecture 1,  $Y$  has Lévy symbol

$$\eta_Y(u) = \left[ \int (e^{i(u,y)} - 1) \lambda \mu_Z(dy) \right].$$

Again the sample paths of  $Y$  are piecewise constant with “jump discontinuities” at the random times  $(T(n), n \in \mathbb{N})$ , however this time the size of the jumps is itself random, and the jump at  $T(n)$  can be any value in the range of the random variable  $Z(n)$ .



Simulation of a compound Poisson process with  $N(0, 1)$

### Example 4 - Interlacing Processes

Let  $C$  be a Gaussian Lévy process as in Example 1 and  $Y$  be a compound Poisson process as in Example 3, which is independent of  $C$ .

Define a new process  $X$  by

$$X(t) = C(t) + Y(t),$$

for all  $t \geq 0$ , then it is not difficult to verify that  $X$  is a Lévy process with Lévy symbol

$$\eta_X(u) = i(b, u) - \frac{1}{2}(u, Au) + \left[ \int (e^{i(u, y)} - 1) \lambda \mu_Z(dy) \right].$$

Using the notation of Examples 2 and 3, we see that the paths of  $X$  have jumps of random size occurring at random times.

$$\begin{aligned} X(t) &= C(t) \quad \text{for } 0 \leq t < T_1, \\ &= C(T_1) + Z_1 \quad \text{when } t = T_1, \\ &= X(T_1) + C(t) - C(T_1) \quad \text{for } T_1 < t < T_2, \\ &= X(T_2-) + Z_2 \quad \text{when } t = T_2, \end{aligned}$$

and so on recursively. We call this procedure an *interlacing* as a continuous path process is “interlaced” with random jumps. It seems reasonable that the most general Lévy process might arise as the limit of a sequence of such interlacings, and this can be established rigorously.

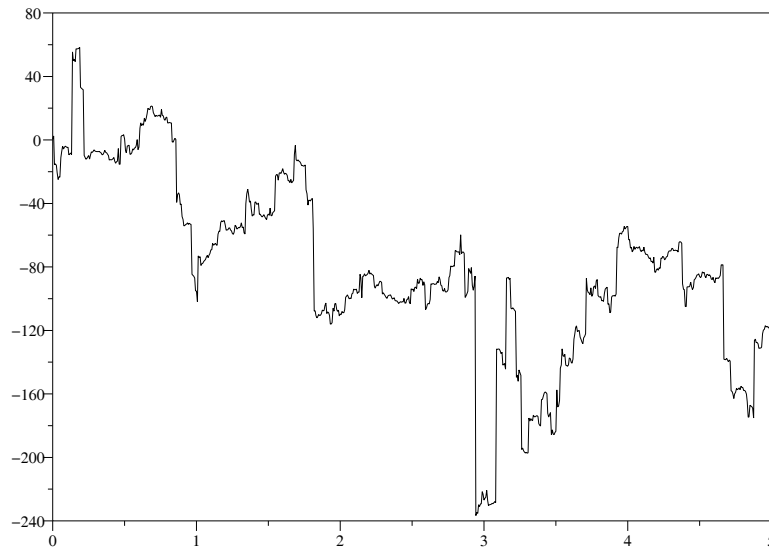
### Example 5 - Stable Lévy Processes

A *stable Lévy process* is a Lévy process  $X$  in which the Lévy symbol is that of a given stable law. So, in particular, each  $X(t)$  is a stable random variable. For example, we have the rotationally invariant case whose Lévy symbol is given by

$$\eta(u) = -\sigma^\alpha |u|^\alpha,$$

where  $\alpha$  is the index of stability ( $0 < \alpha \leq 2$ ). One of the reasons why these are important in applications is that they display self-similarity.

In general, a stochastic process  $Y = (Y(t), t \geq 0)$  is *self-similar with Hurst index*  $H > 0$  if the two processes  $(Y(at), t \geq 0)$  and  $(a^H Y(t), t \geq 0)$  have the same finite-dimensional distributions for all  $a \geq 0$ . By examining characteristic functions, it is easily verified that a rotationally invariant stable Lévy process is self-similar with Hurst index  $H = \frac{1}{\alpha}$ , so that e.g. Brownian motion is self-similar with  $H = \frac{1}{2}$ . A Lévy process  $X$  is self-similar if and only if each  $X(t)$  is strictly stable.



Simulation of the Cauchy process.

Every Lévy process with a non-degenerate Gaussian component has a density.

In this case

$$\Re(\eta(u)) = -\frac{1}{2}(u, Au) + \int_{\mathbb{R}^d - \{0\}} (\cos(u, y) - 1) \nu(dy),$$

and so

$$\int_{\mathbb{R}^d} e^{t\Re(\eta(u))} du \leq \int_{\mathbb{R}^d} e^{-\frac{t}{2}(u, Au)} < \infty,$$

using  $(u, Au) \geq \lambda|u|^2$  where  $\lambda > 0$  is smallest eigenvalue of  $A$ .

## Densities of Lévy Processes

Question: When does a Lévy process have a density  $f_t$  for all  $t > 0$  so that for all Borel sets  $B$ :

$$P(X_t \in B) = p_t(B) = \int_B f_t(x) dx?$$

In general, a random variable has a continuous density if its characteristic function is integrable and in this case, the density is the Fourier transform of the characteristic function.

So for Lévy processes, if for all  $t > 0$ ,

$$\int_{\mathbb{R}^d} |e^{t\eta(u)}| du = \int_{\mathbb{R}^d} e^{t\Re(\eta(u))} du < \infty$$

we then have

$$f_t(x) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{t\eta(u) - i(u, x)} du.$$

For examples where densities exist for  $A = 0$  with  $d = 1$ : if  $X$  is  $\alpha$ -stable, it has a density since for all  $1 \leq \alpha \leq 2$ :

$$\int_{|u| \geq 1} e^{-t|u|^\alpha} du \leq \int_{|u| \geq 1} e^{-t|u|} du < \infty,$$

and for  $0 \leq \alpha < 1$ :

$$\int_{\mathbb{R}} e^{-t|u|^\alpha} du = \frac{2}{\alpha} \int_0^\infty e^{-ty} y^{\frac{1}{\alpha}-1} dy < \infty.$$



In general, a sufficient condition for a density is

- $\nu(\mathbb{R}^d) = \infty$
- $\tilde{\nu}^{*m}$  is absolutely continuous with respect to Lebesgue measure for some  $m \in \mathbb{N}$  where

$$\tilde{\nu}(A) = \int_A (|x|^2 \wedge 1) \nu(dx).$$

A Lévy process has a *Lévy density*  $g_\nu$  if its Lévy measure  $\nu$  is absolutely continuous with respect to Lebesgue measure, then  $g_\nu$  is defined to be the Radon-Nikodym derivative  $\frac{d\nu}{dx}$ .

A process may have a Lévy density but not have a density.

**Example.** Let  $X$  be a compound Poisson process with each  $X(t) = Y_1 + Y_2 + \dots + Y_{N(t)}$  wherein each  $Y_j$  has a density  $f_Y$ , then  $g_\nu = \lambda f_Y$  is the Lévy density.

But

$$P(Y(t) = 0) \geq P(N(t) = 0) = e^{-\lambda t} > 0,$$

so  $p_t$  has an atom at  $\{0\}$ .

We have  $p_t(A) = e^{-\lambda t} \delta_0(A) + \int_A f_t^{ac}(x) dx$ , where for  $x \neq 0$

$$f_t^{ac}(x) = e^{-\lambda t} \sum_{n=1}^{\infty} \frac{(\lambda t)^n}{n!} f_Y^{*n}(x).$$

$f_t^{ac}(x)$  is the conditional density of  $X(t)$  given that it jumps at least once between 0 and  $t$ .

In this case, (2.2) takes the precise form (for  $x \neq 0$ )

$$g_\nu(x) = \lim_{t \downarrow 0} \frac{f_t^{ac}(x)}{t}.$$

## Subordinators

A *subordinator* is a one-dimensional Lévy process which is increasing a.s. Such processes can be thought of as a random model of time evolution, since if  $T = (T(t), t \geq 0)$  is a subordinator we have

$T(t) \geq 0$  for each  $t > 0$  a.s. and  $T(t_1) \leq T(t_2)$  whenever  $t_1 \leq t_2$  a.s.

Now since for  $X(t) \sim N(0, At)$  we have

$P(X(t) \geq 0) = P(X(t) \leq 0) = \frac{1}{2}$ , it is clear that such a process cannot be a subordinator.

## Theorem

If  $T$  is a subordinator then its Lévy symbol takes the form

$$\eta(u) = ibu + \int_{(0,\infty)} (e^{iuy} - 1)\lambda(dy), \quad (2.4)$$

where  $b \geq 0$ , and the Lévy measure  $\lambda$  satisfies the additional requirements

$$\lambda(-\infty, 0) = 0 \text{ and } \int_{(0,\infty)} (y \wedge 1)\lambda(dy) < \infty.$$

Conversely, any mapping from  $\mathbb{R}^d \rightarrow \mathbb{C}$  of the form (2.4) is the Lévy symbol of a subordinator.

We call the pair  $(b, \lambda)$ , the *characteristics* of the subordinator  $T$ .

For each  $t \geq 0$ , the map  $u \rightarrow \mathbb{E}(e^{iuT(t)})$  can be analytically continued to the region  $\{iu, u > 0\}$  and we then obtain the following expression for the Laplace transform of the distribution

$$\mathbb{E}(e^{-uT(t)}) = e^{-t\psi(u)},$$

$$\text{where } \psi(u) = -\eta(iu) = bu + \int_{(0,\infty)} (1 - e^{-uy})\lambda(dy) \quad (2.5)$$

for each  $t, u \geq 0$ .

This is much more useful for both theoretical and practical application than the characteristic function.

The function  $\psi$  is usually called the *Laplace exponent* of the subordinator.

## Examples of Subordinators

### (1) The Poisson Case

Poisson processes are clearly subordinators. More generally a compound Poisson process will be a subordinator if and only if the  $Z(n)$ 's are all  $\mathbb{R}^+$  valued.

### (2) $\alpha$ -Stable Subordinators

Using straightforward calculus, we find that for  $0 < \alpha < 1$ ,  $u \geq 0$ ,

$$u^\alpha = \frac{\alpha}{\Gamma(1-\alpha)} \int_0^\infty (1 - e^{-ux}) \frac{dx}{x^{1+\alpha}}.$$

Hence for each  $0 < \alpha < 1$  there exists an  $\alpha$ -stable subordinator  $T$  with Laplace exponent

$$\psi(u) = u^\alpha.$$

and the characteristics of  $T$  are  $(0, \lambda)$  where  $\lambda(dx) = \frac{\alpha}{\Gamma(1-\alpha)} \frac{dx}{x^{1+\alpha}}$ .

Note that when we analytically continue this to obtain the Lévy symbol we obtain the form given in Lecture 1 for stable laws with  $\mu = 0, \beta = 1$  and  $\sigma^\alpha = \cos\left(\frac{\alpha\pi}{2}\right)$ .

### (3) The Lévy Subordinator

The  $\frac{1}{2}$ -stable subordinator has a density given by the Lévy distribution (with  $\mu = 0$  and  $\sigma = \frac{t^2}{2}$ )

$$f_{T(t)}(s) = \left(\frac{t}{2\sqrt{\pi}}\right) s^{-\frac{3}{2}} e^{-\frac{t^2}{4s}},$$

for  $s \geq 0$ . The Lévy subordinator has a nice probabilistic interpretation as a first hitting time for one-dimensional standard Brownian motion ( $B(t), t \geq 0$ ),

$$T(t) = \inf \left\{ s > 0; B(s) = \frac{t}{\sqrt{2}} \right\}. \quad (2.6)$$

To show directly that for each  $t \geq 0$ ,

$$\mathbb{E}(e^{-uT(t)}) = \int_0^\infty e^{-us} f_{T(t)}(s) ds = e^{-tu^{\frac{1}{2}}},$$

write  $g_t(u) = \mathbb{E}(e^{-uT(t)})$ . Differentiate with respect to  $u$  and make the substitution  $x = \frac{t^2}{4us}$  to obtain the differential equation  $g_t'(u) = -\frac{t}{2\sqrt{u}} g_t(u)$ . Via the substitution  $y = \frac{t}{2\sqrt{s}}$  we see that  $g_t(0) = 1$  and the result follows.

### (4) Inverse Gaussian Subordinators

We generalise the Lévy subordinator by replacing Brownian motion by the Gaussian process  $C = (C(t), t \geq 0)$  where each  $C(t) = B(t) + \mu t$  and  $\mu \in \mathbb{R}$ . The *inverse Gaussian subordinator* is defined by

$$T(t) = \inf\{s > 0; C(s) = \delta t\}$$

where  $\delta > 0$  and is so-called since  $t \rightarrow T(t)$  is the generalised inverse of a Gaussian process.

Using martingale methods, we can show that for each  $t, u > 0$ ,

$$\mathbb{E}(e^{-uT(t)}) = e^{-t\delta(\sqrt{2u+\mu^2}-\mu)}, \quad (2.7)$$

In fact each  $T(t)$  has a density:-

$$f_{T(t)}(s) = \frac{\delta t}{\sqrt{2\pi}} e^{\delta t \mu} s^{-\frac{3}{2}} \exp \left\{ -\frac{1}{2}(t^2 \delta^2 s^{-1} + \mu^2 s) \right\}, \quad (2.8)$$

for each  $s, t \geq 0$ .

In general any random variable with density  $f_{T(t)}$  is called an *inverse Gaussian* and denoted as  $IG(\delta, \mu)$ .

### (5) Gamma Subordinators

Let  $(T(t), t \geq 0)$  be a *gamma process* with parameters  $a, b > 0$  so that each  $T(t)$  has density

$$f_{T(t)}(x) = \frac{b^{at}}{\Gamma(at)} x^{at-1} e^{-bx},$$

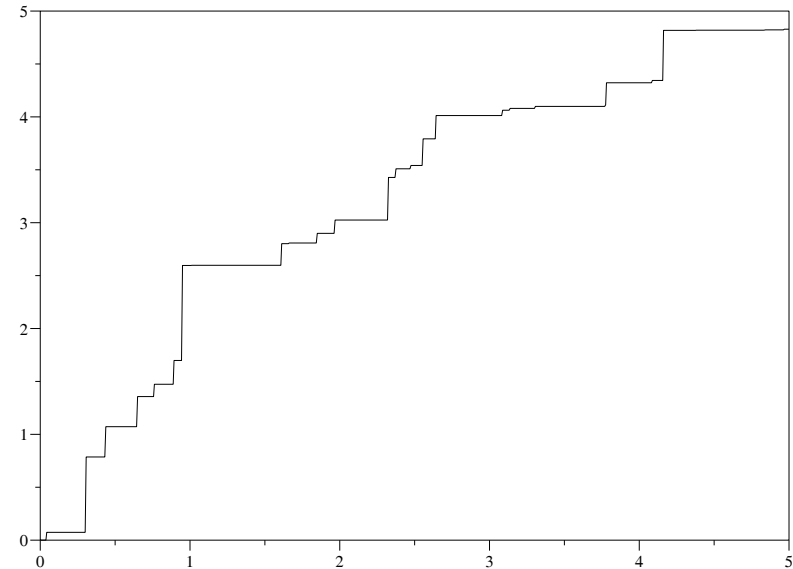
for  $x \geq 0$ ; then it is easy to verify that for each  $u \geq 0$ ,

$$\int_0^\infty e^{-ux} f_{T(t)}(x) dx = \left(1 + \frac{u}{b}\right)^{-at} = \exp\left(-t a \log\left(1 + \frac{u}{b}\right)\right).$$

From here it is a straightforward exercise in calculus to show that

$$\int_0^\infty e^{-ux} f_{T(t)}(x) dx = \exp\left[-t \int_0^\infty (1 - e^{-ux}) a x^{-1} e^{-bx} dx\right].$$

From this we see that  $(T(t), t \geq 0)$  is a subordinator with  $b = 0$  and  $\lambda(dx) = ax^{-1}e^{-bx}dx$ . Moreover  $\psi(u) = a \log(1 + \frac{u}{b})$  is the associated Bernstein function (see below).



Simulation of a gamma subordinator.

Before we go further into the probabilistic properties of subordinators we'll make a quick diversion into analysis. Let  $f \in C^\infty((0, \infty))$ . We say it is *completely monotone* if  $(-1)^n f^{(n)} \geq 0$  for all  $n \in \mathbb{N}$ , and a *Bernstein function* if  $f \geq 0$  and  $(-1)^n f^{(n)} \leq 0$  for all  $n \in \mathbb{N}$ .

### Theorem

- 1  $f$  is a Bernstein function if and only if the mapping  $x \rightarrow e^{-tf(x)}$  is completely monotone for all  $t \geq 0$ .
- 2  $f$  is a Bernstein function if and only if it has the representation

$$f(x) = a + bx + \int_0^\infty (1 - e^{-yx})\lambda(dy),$$

for all  $x > 0$  where  $a, b \geq 0$  and  $\int_0^\infty (y \wedge 1)\lambda(dy) < \infty$ .

- 3  $g$  is completely monotone if and only if there exists a measure  $\mu$  on  $[0, \infty)$  for which

$$g(x) = \int_0^\infty e^{-xy}\mu(dy).$$

To interpret this theorem, first consider the case  $a = 0$ . In this case, if we compare the statement of Theorem 4 with equation (2.5), we see that there is a one to one correspondence between Bernstein functions for which  $\lim_{x \rightarrow 0} f(x) = 0$  and Laplace exponents of subordinators. The Laplace transforms of the laws of subordinators are always completely monotone functions and a subclass of all possible measures  $\mu$  appearing in Theorem 4 (3) is given by all possible laws  $p_{T(t)}$  associated to subordinators. A general Bernstein function with  $a > 0$  can be given a probabilistic interpretation by means of “killing”.

One of the most important probabilistic applications of subordinators is to “time change”. Let  $X$  be an arbitrary Lévy process and let  $T$  be a subordinator defined on the same probability space as  $X$  such that  $X$  and  $T$  are independent. We define a new stochastic process  $Z = (Z(t), t \geq 0)$  by the prescription

$$Z(t) = X(T(t)),$$

for each  $t \geq 0$  so that for each  $\omega \in \Omega$ ,  $Z(t)(\omega) = X(T(t)(\omega))(\omega)$ . The key result is then the following.

**Theorem**  
*Z is a Lévy process.*

We compute the Lévy symbol of the subordinated process  $Z$ .

**Theorem**  

$$\eta_Z = -\psi_T \circ (-\eta_X).$$

*Proof.* For each  $u \in \mathbb{R}^d, t \geq 0$ ,

$$\begin{aligned} e^{i\eta_Z(t)(u)} &= \mathbb{E}(e^{i(u, Z(t))}) \\ &= \mathbb{E}(e^{i(u, X(T(t)))}) \\ &= \int \mathbb{E}(e^{i(u, X(s))}) p_{T(t)}(ds) \\ &= \int e^{s\eta_X(u)} p_{T(t)}(ds) \\ &= \mathbb{E}(e^{-(-\eta_X(u))T(t)}) \\ &= e^{-t\psi_T(-\eta_X(u))}. \quad \square \end{aligned}$$

**Example : From Brownian Motion to  $2\alpha$ -stable Processes**

Let  $T$  be an  $\alpha$ -stable subordinator (with  $0 < \alpha < 1$ ) and  $X$  be a  $d$ -dimensional Brownian motion with covariance  $A = 2I$ , which is independent of  $T$ . Then for each  $s \geq 0, u \in \mathbb{R}^d, \psi_T(s) = s^\alpha$  and  $\eta_X(u) = -|u|^2$ , and hence  $\eta_Z(u) = -|u|^{2\alpha}$ , i.e.  $Z$  is a rotationally invariant  $2\alpha$ -stable process.

In particular, if  $d = 1$  and  $T$  is the Lévy subordinator, then  $Z$  is the *Cauchy process*, so each  $Z(t)$  has a symmetric Cauchy distribution with parameters  $\mu = 0$  and  $\sigma = 1$ . It is interesting to observe from (2.6) that  $Z$  is constructed from two independent standard Brownian motions.

Examples of subordinated processes have recently found useful applications in mathematical finance. We briefly mention two interesting cases:-

(i) **The Variance Gamma Process**

In this case  $Z(t) = B(T(t))$ , for each  $t \geq 0$ , where  $B$  is a standard Brownian motion and  $T$  is an independent gamma subordinator. The name derives from the fact that, in a formal sense, each  $Z(t)$  arises by replacing the variance of a normal random variable by a gamma random variable. Using Theorem 6, a simple calculation yields

$$\Phi_{Z(t)}(u) = \left(1 + \frac{u^2}{2b}\right)^{-at},$$

for each  $t \geq 0, u \in \mathbb{R}$ , where  $a$  and  $b$  are the usual parameters which determine the gamma process. It is an easy exercise in manipulating characteristic functions to compute the alternative representation:

$$Z(t) = G(t) - L(t),$$

where  $G$  and  $L$  are independent gamma subordinators each with parameters  $\sqrt{2b}$  and  $a$ . This yields a nice financial representation of  $Z$  as a difference of independent “gains” and “losses”. From this representation, we can compute that  $Z$  has a Lévy density

$$g_\nu(x) = \frac{a}{|x|} (e^{\sqrt{2b}x} \mathbf{1}_{(-\infty,0)}(x) + e^{-\sqrt{2b}x} \mathbf{1}_{(0,\infty)}(x)).$$

The *CGMY processes* are a generalisation of the variance-gamma processes due to Carr, Geman, Madan and Yor. They are characterised by their Lévy density:

$$g_\nu(x) = \frac{a}{|x|^{1+\alpha}} (e^{b_1 x} \mathbf{1}_{(-\infty,0)}(x) + e^{-b_2 x} \mathbf{1}_{(0,\infty)}(x)),$$

where  $a > 0, 0 \leq \alpha < 2$  and  $b_1, b_2 \geq 0$ . We obtain stable Lévy processes when  $b_1 = b_2 = 0$ . They can also be obtained by subordinating Brownian motion with drift. The CGMY processes are a subclass of the *tempered stable processes*. Note how the exponential dampens the effects of large jumps.

(ii) **The Normal Inverse Gaussian Process**

In this case  $Z(t) = C(T(t)) + \mu t$  for each  $t \geq 0$  where each  $C(t) = B(t) + \beta t$ , with  $\beta \in \mathbb{R}$ . Here  $T$  is an inverse Gaussian subordinator, which is independent of  $B$ , and in which we write the parameter  $\gamma = \sqrt{\alpha^2 - \beta^2}$ , where  $\alpha \in \mathbb{R}$  with  $\alpha^2 \geq \beta^2$ .  $Z$  depends on four parameters and has characteristic function

$$\Phi_{Z(t)}(\alpha, \beta, \delta, \mu)(u) = \exp \{ \delta t (\sqrt{\alpha^2 - \beta^2} - \sqrt{\alpha^2 - (\beta + iu)^2}) + i\mu t u \}$$

for each  $u \in \mathbb{R}, t \geq 0$ . Here  $\delta > 0$  is as in (2.7). Each  $Z(t)$  has a density given by

$$f_{Z(t)}(x) = C(\alpha, \beta, \delta, \mu; t) q \left( \frac{x - \mu t}{\delta t} \right)^{-1} K_1 \left( \delta t \alpha q \left( \frac{x - \mu t}{\delta t} \right) \right) e^{\beta x},$$

for each  $x \in \mathbb{R}$ , where

$q(x) = \sqrt{1 + x^2}$ ,  $C(\alpha, \beta, \delta, \mu; t) = \pi^{-1} \alpha e^{\delta t \sqrt{\alpha^2 - \beta^2} - \beta \mu t}$  and  $K_1$  is a Bessel function of the third kind.